The Effect of Ionic Radius Difference of RE^{3+} and Ba^{2+} on the Shock Synthesis of $REBa_2Cu_3O_{\nu}$ (RE = Y, Eu, La)

Hideaki Hikosaka, Keiji Kusaba, Kiyoto Fukuoka, Masae Kikuchi, and Yasuhiko Syono

Institute for Materials Research, Tohoku University, Sendai 980-77, Japan

Received April 7, 1997; in revised form September 22, 1997; accepted October 12, 1997

Shock-synthesis experiments of $REBa_2Cu_3O_{\nu}$ (RE = Y, Eu, and La) were carried out for the mixture of component oxides as starting materials using a 25 mm-bore propellant gun. Both tetragonal EuBa₂Cu₃O_y with c/a = 3.00 and cubic LaBa₂Cu₃O_y were successfully synthesized by shock compression, but YBa₂ Cu₃O_y was synthesized only partly. The optimum condition for obtaining maximum yield of shock-synthesized EuBa₂Cu₃O_v and LaBa₂Cu₃O_v was shock compression of the starting materials with 40% porosity to 30 GPa. $La_{1+x}Ba_{2-x}$ Cu₃O_y could also be synthesized, but Eu_{1+x}Ba_{2-x}Cu₃O_y could not. Shock-synthesized EuBa₂Cu₃O_v and LaBa₂Cu₃O_v were in a metastable phase in which RE³⁺ and Ba²⁺ were not ordered perfectly; they converted to the phase synthesized by conventional solid state reaction when annealed above 900°C. Only after annealing above 900°C, followed by oxidation at 400°C, did EuBa₂Cu₃O₂ exhibit superconductivity with $T_c = 71$ K. These results suggest that shock synthesis of REBa₂Cu₃O_v becomes easier with a decrease in the difference between ionic radii of RE³⁺ and Ba²⁺. © 1998 Academic Press

INTRODUCTION

The shock compression technique has been noted as a unique method for synthesizing materials within a very short time. In the shock compression of a powder specimen, rapid chemical reaction is either directly induced or assisted by the introduction of a large concentration of defect. The residual temperature after pressure release remains for only several milliseconds, and metastable phase can be synthesized because of the highly nonequilibrium process. Shock compression method has been successfully applied to the synthesis of high-temperature oxide superconductors such as $La_{2-x}Sr_xCuO_4$ (1, 2) and $Tl_2Ba_2CuO_6$ (3). We already reported the shock synthesis of EuBa₂Cu₃O_v, which was found to be a metastable tetragonal phase with c/a = 3.00, in contrast to the high T_c phase synthesized by conventional solid state reaction (4). This difference has been explained as incomplete ordering among the metal ions of the A site because of the short duration of the shock process. $Eu_{1+x}Ba_{2-x}Cu_{3}O_{y}$ with $0.1 \le x \le 0.5$ could not be synthesized by shock compression.

In the present study, we have synthesized REBa₂Cu₃O_y (RE = Y, Eu, and La) by shock compression and compared the effect of ionic radius of RE³⁺ on synthesis conditions and resulting crystal structure. Shock synthesis of RE_{1+x} Ba_{2-x}Cu₃O_y was also carried out for RE = Eu³⁺ and La³⁺, where a wide solid solution range for $0 \le x \le 0.5$ exists (5). Furthermore, the effect of annealing shock-synthesized EuBa₂Cu₃O_y and LaBa₂Cu₃O_y was examined.

EXPERIMENTAL

Powder materials of component oxide, RE_2O_3 (RE = Y, Eu, and La), CuO, and BaCuO₂ were mixed in a stoichiometric ratio by ball milling for 3 h to give $RE_{1+x}Ba_{2-x}Cu_3$ O_{v} . The mixture was pressed into a pellet 13 mm in diameter and 1 mm in thickness. The pellet was encapsulated in a SUS304 stainless steel container. Shock loading experiments were carried out using a powder gun with 25 mm bore (6). The specimen assembly was impacted by a SUS304 stainless steel flyer with 3 mm thickness in the flyer velocity range 1.0–1.7 km/s. The shock pressure in the container was calculated to be 20-40 GPa from the measured flyer velocity on the basis of the impedance match concept (7). The shock pressure achieved in the specimen container was taken to represent the shocked state of the sample by assuming equilibration with the container via multiple reflection across the specimen-container boundaries. The shock-synthesis conditions were controlled by the shock pressure and porosity of the starting materials in the range of 50-70% of theoretical density.

The recovered specimens were examined by X-ray powder diffraction (XRD) analysis with $CuK\alpha$ radiation. The oxygen content was determined by iodometric titration. Magnetic susceptibility measurements were carried out using a SQUID magnetometer.

RESULTS AND DISCUSSION

X-ray powder diffraction analysis of the shock-recovered specimen of $REBa_2Cu_3O_y$ showed the formation of a

perovskite-based phase. The shock conditions, pressure and porosity, were changed over a wide range to identify the best shock-synthesis conditions for REBa₂Cu₃O_y. It was considered that the reaction took place with shock-induced temperature increases when powder material was shockloaded. Figure 1 shows the relationship between apparent yield of shock-synthesized REBa₂Cu₃O_y and shock temperature, which was estimated from shock pressure and porosity of starting materials by assuming the Hugoniot of $BaTiO_3$ (8) with the perovskite structure. The small yield of $YBa_2Cu_3O_{\nu}$ hardly changed when the shock temperature was varied. The shock-synthesis yield of $YBa_2Cu_3O_{\nu}$ did not increase above 17%. The yield of both $EuBa_2Cu_3O_{\nu}$ and LaBa₂Cu₃O_y increased with increasing shock temperature above 1000°C, and maximum yield was obtained at around 1500°C. There was no appreciable difference between optimum shock temperature for the synthesis of EuBa₂Cu₃O_y and LaBa₂Cu₃O_y. The best shock-synthesis conditions were achieved by shock compression to 30 GPa for a 60% packing density pellet, which corresponds to shock temperature of around 1500°C.

Figure 2 shows XRD patterns of shock-recovered specimens of $REBa_2Cu_3O_y$ with the best shock conditions. In the XRD pattern of $YBa_2Cu_3O_y$, diffraction lines due to a triple-perovskite phase were observed, indicating the synthesis of $YBa_2Cu_3O_y$, although a considerable portion of



FIG. 1. Yield of shock-synthesized REBa₂Cu₃O_y plotted as a function of shock temperature. Triangles, circles, and squares indicate $YBa_2Cu_3O_y$, $EuBa_2Cu_3O_y$, and $LaBa_2Cu_3O_y$, respectively.



FIG. 2. X-ray powder diffraction patterns of REBa₂Cu₃O_y shock-synthesized at 30 GPa: (a) YBa₂Cu₃O_y (closed circles and solid triangles indicate raw materials and BaCO₃); (b) EuBa₂Cu₃O_y (open triangles indicate Eu₂CuO₄); and (c) LaBa₂Cu₃O_y (open circles indicate La₂CuO₄).

unreacted starting materials remained and BaCO₃ was formed. We considered that BaCO₃ was formed after the shock-loaded sample was taken out of the capsule. BaCuO₂ was partially decomposed to BaO and CuO when shockloaded. BaO was not stable in air; therefore the reaction $BaO + CO_2 \rightarrow BaCO_3$ could occur. The shock-synthesized phase was indexed with tetragonal indices with c/a = 3.00, which is smaller than that of $YBa_2Cu_3O_{y}$ synthesized by conventional solid state reaction. The XRD pattern of shock-synthesized EuBa₂Cu₃O_v indicated an almost single phase of triple-perovskite EuBa₂Cu₃O_v, which was also indexed with a tetragonal cell with c/a = 3.00. Cell constants of shock-synthesized EuBa₂Cu₃O_v were smaller than those of the high T_c phase synthesized by conventional solid state reaction. In the case of La³⁺, which has the largest ionic radius among rare earth ions, an almost single phase of $LaBa_2Cu_3O_{\nu}$ was synthesized, as in the case of EuBa₂Cu₃O_y. However, the XRD pattern of shock-synthesized LaBa₂Cu₃O_y was indexed with a cubic cell, and a period three times that of the basic perovskite unit was not observed in the electron diffraction pattern in LaBa₂Cu₃O_v,

unlike $EuBa_2Cu_3O_y$. These results suggest that all shocksynthesized phases are metastable phase whose metal ions of the A site are not ordered perfectly, because of the short duration of the shock process. Shock synthesis of $REBa_2Cu_3O_y$ became easier with a decrease in the difference between the ionic radii of RE^{3+} and Ba^{2+} because of easier disordering.

There was not much difference in shock synthesis of EuBa₂Cu₃O_v and LaBa₂Cu₃O_v. However, a remarkable difference between RE = La and Eu was observed in the case of shock synthesis of RE_{1.5}Ba_{1.5}Cu₃O_v, as shown in the XRD patterns of the recovered specimen (Fig. 3). A large amount of Eu₂CuO₄ was found in the shock-loaded product, and no $Eu_{1,5}Ba_{1,5}Cu_3O_{\nu}$ phase was formed. On the other hand, La_{1.5}Ba_{1.5}Cu₃O_v was successfully synthesized as an almost single phase by shock compression and its XRD pattern was indexed with a cubic cell, the same as LaBa₂Cu₃O_y. Furthermore La_{1+x}Ba_{2-x}Cu₃O_y was found to be synthesized for x = 0.1, 0.3, and 0.4. This result can be explained by the smaller differences between the ionic radii of $La^{3+}(1.18 \text{ Å})$ and $Ba^{2+}(1.42 \text{ Å})$ than between those of $Eu^{3+}(1.07 \text{ Å})$ and Ba^{2+} . The cubic cell constants of all the shock-synthesized $La_{1+x}Ba_{2-x}Cu_3O_y$ are plotted against La³⁺ content in Fig. 4. The cell constants of both shocksynthesized $La_{1+x}Ba_{2-x}Cu_3O_y$ and the specimens synthesized by conventional solid state reaction decreased with increasing La³⁺ content. However, the cell constants of the former were considerably smaller than those of the latter. This result suggests that all shock-synthesized



FIG. 3. X-ray powder diffraction patterns of $RE_{1.5}Ba_{1.5}Cu_3O_y$ shocksynthesized at 30 GPa: (a) $Eu_{1.5}Ba_{1.5}Cu_3O_y$ (closed circles, solid triangles, and open triangles indicate raw materials, $BaCO_3$, and Eu_2CuO_4); and (b) $La_{1.5}Ba_{1.5}Cu_3O_y$ (open circles indicate La_2CuO_4).



FIG. 4. Cell constant of $La_{1+x}Ba_{2-x}Cu_3O_y$ plotted as a function of La^{3+} content. Closed circles indicate *a*-axis of shock-synthesized specimen. Open circles and solid triangles indicate *a*-, *b*-axes and *c*-axis of the specimen synthesized by conventional solid state reaction.

 $La_{1+x}Ba_{2-x}Cu_3O_y$ are metastable phase, probably due to incomplete order among La^{3+} and Ba^{2+} .

The shock-synthesized phases were not superconductors, and an annealing treatment was necessary to achieve superconductivity. The annealing experiments of shock-synthesized EuBa₂Cu₃O_v and LaBa₂Cu₃O_v were carried out. Figure 5 shows the change of the c/a ratio of shock-synthesized EuBa₂Cu₃O_v plotted as a function of annealing temperature in O₂ for 2 h. The c/a ratio of EuBa₂Cu₃O_y remained at 3.00 for annealing below 900°C. Above 900°C the c/a ratio began to increase and reached 3.05, which is the c/a value of EuBa₂Cu₃O_y synthesized by conventional method. Annealing of shock-synthesized EuBa₂Cu₃O_y at 400° C in O₂ did not indicate superconductivity, although the oxygen content increased to 6.77 and Cu valence increased to 2.18, which should show superconductivity. However, the shock-synthesized specimen annealed above 900°C exhibited superconductivity with $T_c = 71$ K (Fig. 6) if its oxygen content was increased to 6.77 by annealing at 400°C in O₂. Shock-synthesized EuBa₂Cu₃O_y consisted of domains whose *c*-axes were orthogonal to each other from HRTEM image observation (4). These results suggest that the complete CuO₂ plane has not been achieved in the shock-synthesized EuBa₂Cu₃O_y because of oxygen disordering originating from a partial disorder between Eu³⁺ and Ba²⁺, and the perfect ordering necessary for achieving



FIG. 5. The c/a ratio of shocked EuBa₂Cu₃O_y annealed in O₂ for 2 h plotted as a function of annealing temperature.

superconductivity is considered to be established above 900° C, where cation migration takes place.

Figure 7 shows the change of cell constant of shocksynthesized LaBa₂Cu₃O_y plotted as a function of annealing temperature in O₂. The cell constant of LaBa₂Cu₃O_y began to increase from 800°C. The crystal system changed from cubic to tetragonal above 900°C. These results also suggest that La³⁺ and Ba²⁺ began to order above 900°C. However, the c/a ratio remained 3.00 even with annealing at 1050°C, indicating that the ordering between La³⁺ and Ba²⁺ was less complete than between Eu³⁺ and Ba²⁺ because of smaller differences in ionic radii.



FIG. 6. Temperature dependence of DC susceptibility of EuBa₂ Cu₃O_{6.77} shock-synthesized and annealed at 900°C, for zero-field cooling (closed circles) and field cooling (open circles) (H = 20 Oe).



FIG. 7. The cell constant of shocked $LaBa_2Cu_3O_y$ annealed in O_2 for 2 h plotted as a function of annealing temperature.

CONCLUSION

Tetragonal EuBa₂Cu₃O_y and cubic La_{1+x}Ba_{2-x}Cu₃O_y were successfully synthesized as an almost single phase by shock compression, but YBa₂Cu₃O_v was synthesized only partly. Shock-synthesized phases seemed to be a metastable phase whose RE^{3+} and Ba^{2+} were not perfectly ordered, and the degree of disordering increased with a decrease in the difference between ionic radii of RE^{3+} and Ba^{2+} . The shock-synthesized phase was considered to have an incomplete CuO₂ plane because of incomplete ordering between RE^{3+} and Ba^{2+} , and did not show superconductivity simply by oxygen treatment. Only after annealing above 900°C did the shock-synthesized EuBa₂Cu₃O_v and LaBa₂Cu₃O_v become superconductors, because RE^{3+} and Ba^{2+} began to order completely. Shock synthesis of REBa₂Cu₃O_y was found to be easier with the decrease in the difference between ionic radii of RE³⁺ and Ba²⁺, because disordering became easier, which favored synthesis due to the short reaction time in the shock process.

ACKNOWLEDGMENT

The authors are grateful to Dr. T. Atou, IMR, Tohoku University for useful discussion.

REFERENCES

1. R. A. Graham, E. L. Venturini, B. Morosin, and D. S. Ginley, *Phys. Lett.* A **123**, 87 (1987).

- B. Morosin, R. A. Graham, E. L. Venturini, D. S. Ginley, and W. F. Hammetter, *in* "Shock Waves in Condensed Matter 1987" (S. C. Schmidt and N. C. Holmes, Eds.), p. 439. Elsevier, Amsterdam, 1988.
- Z. Iqbal, N. N. Thadhani, N. Chawla, B. L. Ramakrishna, R. Sharma, K. V. Rao, S. Skumryev, F. Reidinger, and H. Eckhardt, *Appl. Phys. Lett.* 55, 2339 (1989).
- 4. H. Hikosaka, T. Atou, K. Kusaba, T. Suzuki, K. Fukuoka, M. Kikuchi, and Y. Syono, Jpn. J. Appl. Phys. 34, 1507 (1995).
- W. Wong-Ng, B. Paretzkin, and E. R. Fuller Jr., J. Solid State Chem. 85, 117 (1990).
- 6. T. Goto and Y. Syono, in "Materials Science of Earth's Interior" (I. Sunagawa, Ed.), p. 605. Terra, Tokyo, 1984.
- R. G. McQueen, S. P. Marsh, J. W. Taylor, J. N. Fritz, and W. J. Carter, in "High-Velocity Impact Phenomena" (R. Kinslow, Ed.), p. 293. Academic Press, New York, 1970.
- S. P. Marsh (Ed.), "LASL Shock Hugoniot Data," p. 240. University of California Press, Berkeley, 1980.